

primordial magnetic fields during the electroweak crossover and baryogenesis from magnetic helicity decay

[2012.14435], [2507.01576], [2509.23858], [2509.25734]

2026/2/12, Tohoku U.

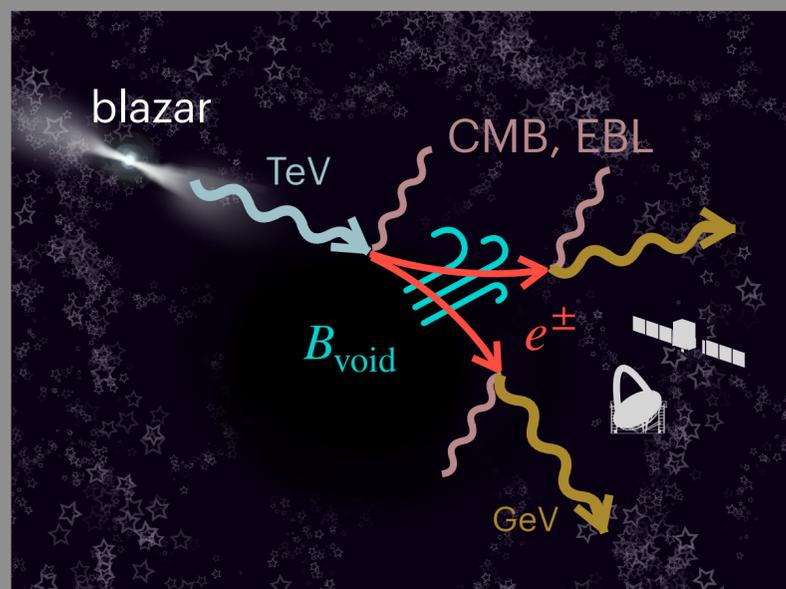
Fumio Uchida (CTPU-CGA, IBS)

Big questions!

magnetogenesis

&

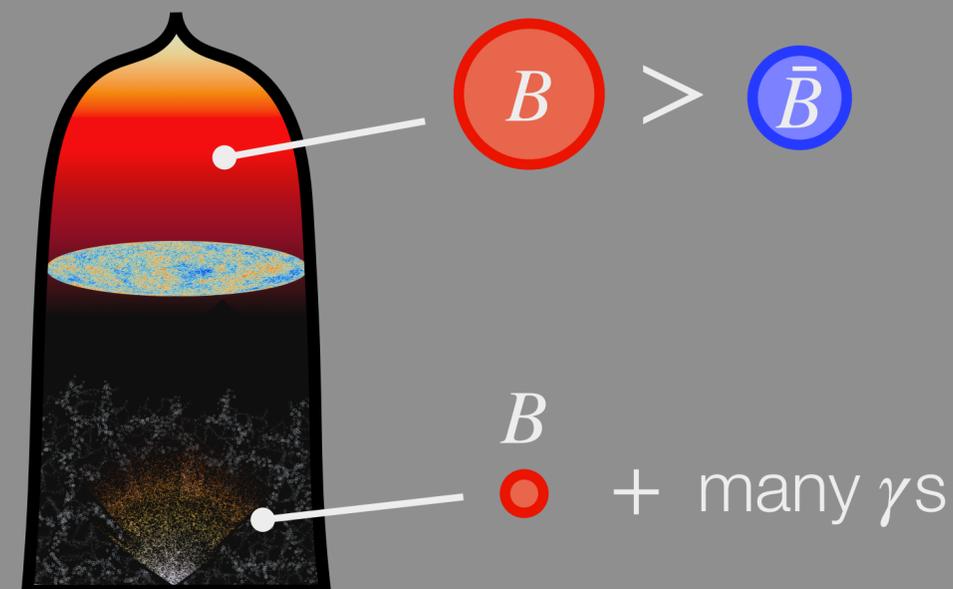
baryogenesis



$$B_{\text{void}} \gtrsim 10^{-17} \text{ G}$$

[Neronov, Vovk 2010], ...

How to generate \vec{B} initially? $B_{\text{void}} \sim B_{\text{PMF}}$?

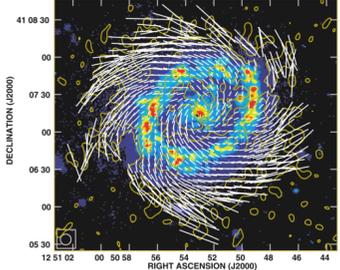


$$n_B \sim 10^{-9} n_\gamma$$

[Planck 2018], ...

How to generate $n_B - n_{\bar{B}}$ initially?

Magnetogenesis



[Chyży, Buta]

observed in stars, planets, ...
 observed in galaxies $\sim 10^{-6}$ G at ~ 10 kpc

sizable B

astrophysics
 structure formation

seed magnetic field $\gtrsim 10^{-20}$ G at $\gtrsim 10$ kpc (?)

tiny B

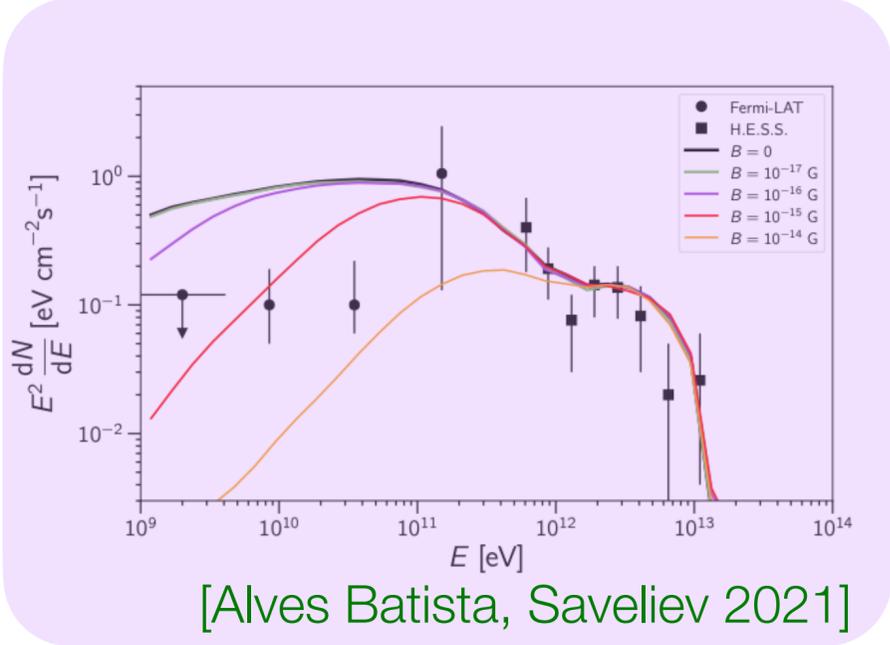
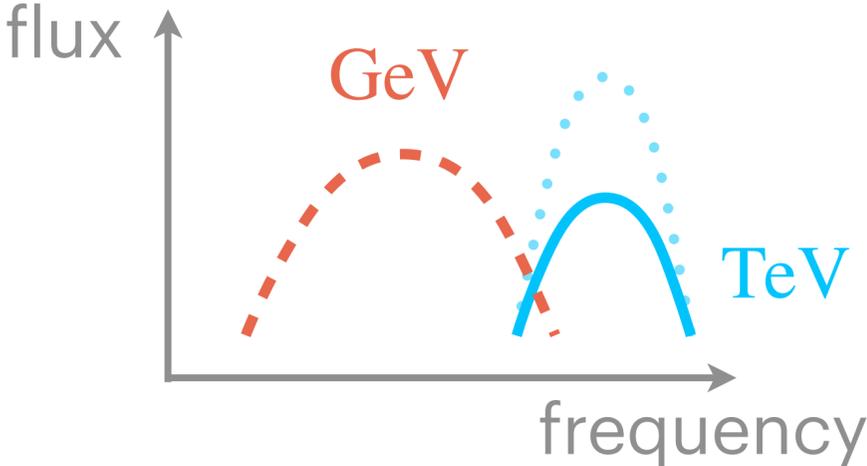
“magnetogenesis”
 involves BSM (?)

vacuum fluctuation $\sim 10^{-50-60}$ G at 1 Mpc

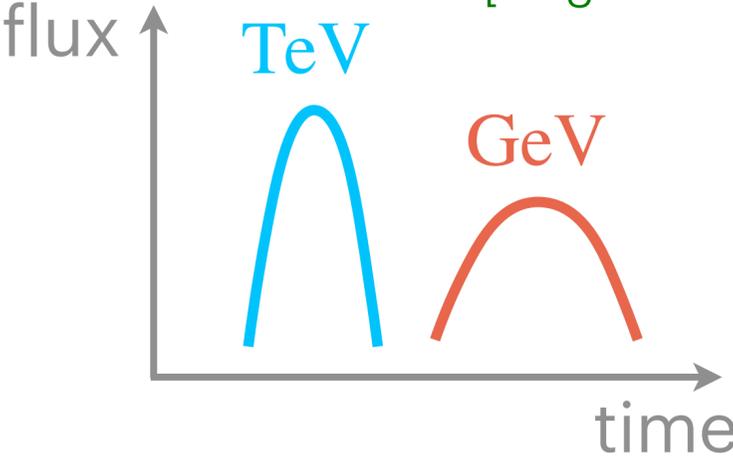
$B \simeq 0$

Void magnetic field

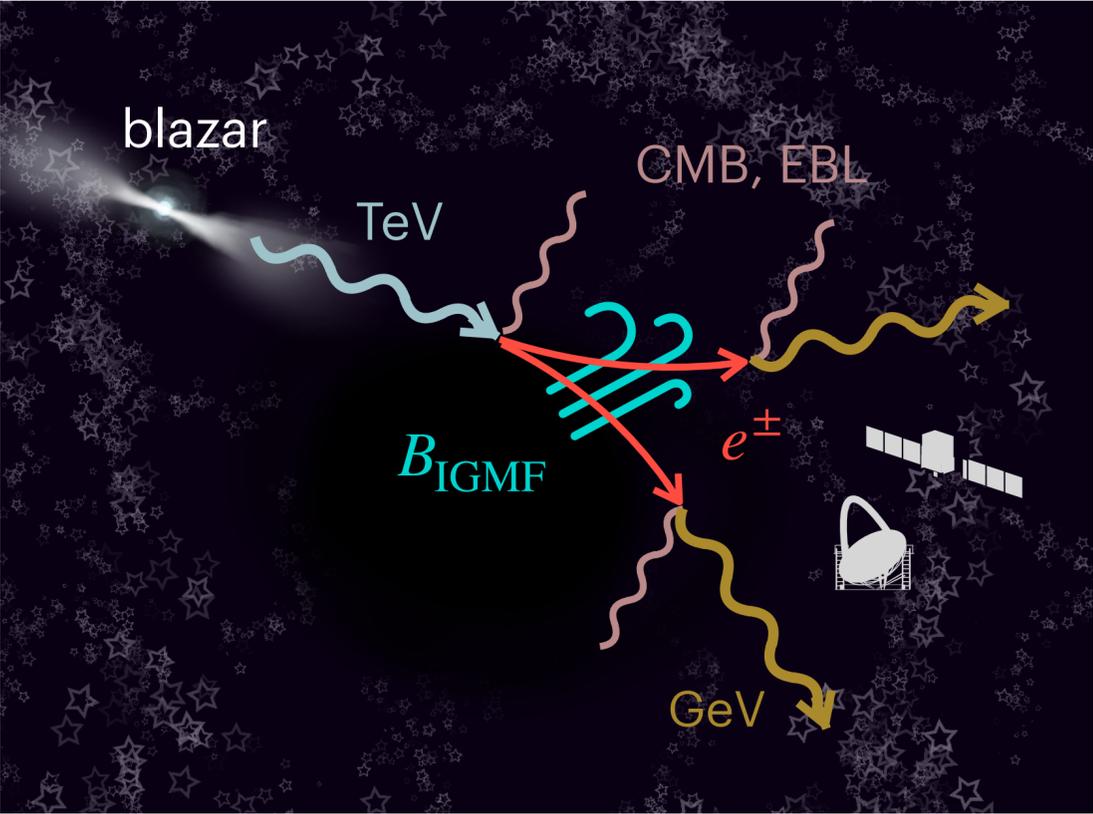
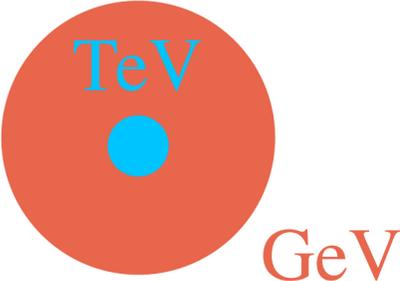
- lack of GeV photons
[Honda 1989]



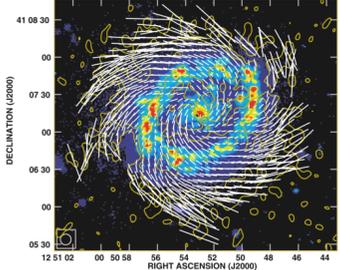
- effects on GeV photons, if observed
- time delay (pair echo) [Plaga 1995]



- diffuse flux (pair halo)
[Honda 1989] [Aharonian+ 1994]
[Plaga 1995]



Magnetogenesis



[Chyży, Buta]

observed in stars, planets, ...
 observed in galaxies $\sim 10^{-6}$ G at ~ 10 kpc

sizable B

astrophysics
 structure formation

observed in voids $\gtrsim 10^{-17}$ G

↕ equivalent (?)

seed magnetic field $\gtrsim 10^{-20}$ G at $\gtrsim 10$ kpc (?)

feedbacks (?)

tiny B

“magnetogenesis”
 involves BSM (?)

vacuum fluctuation $\sim 10^{-50-60}$ G at 1 Mpc

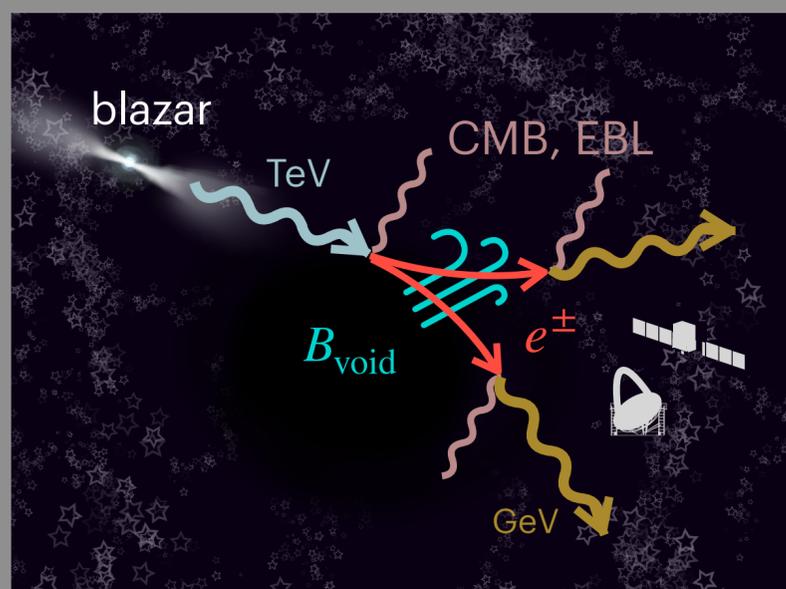
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Big questions!

magnetogenesis

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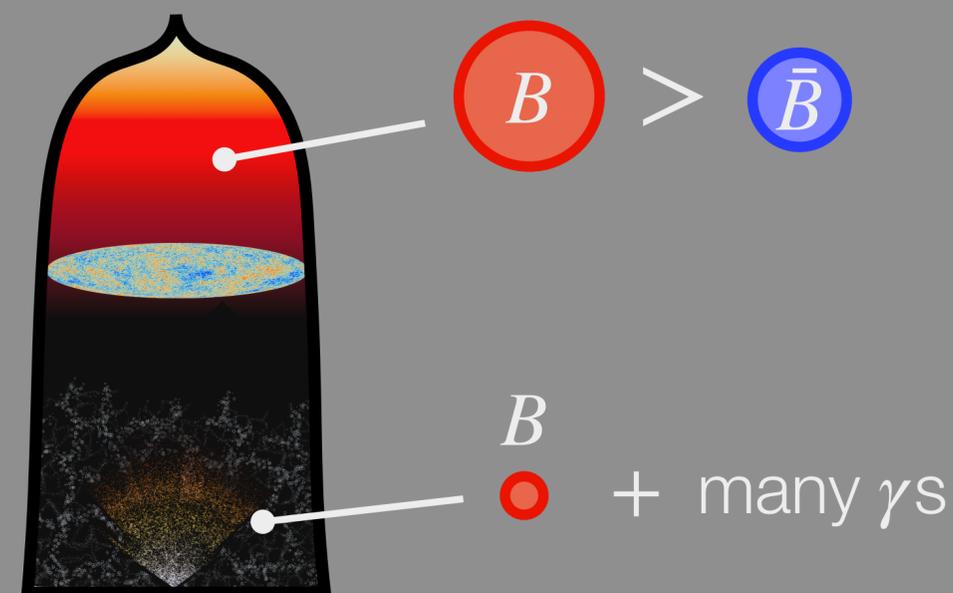
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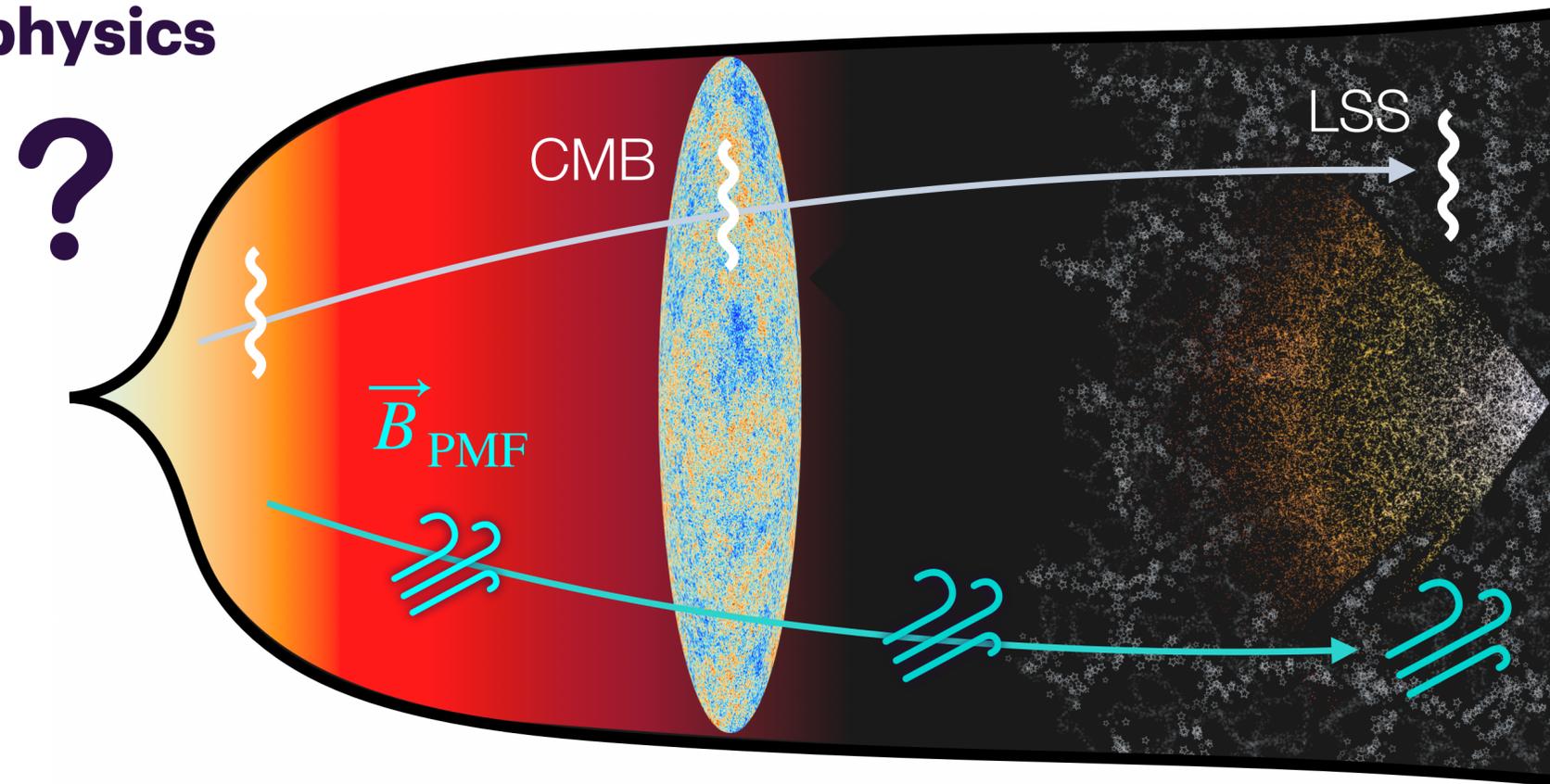
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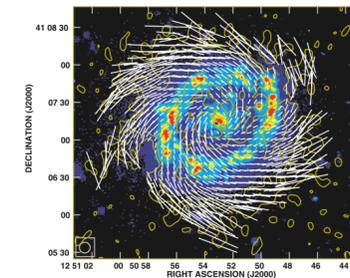
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This talk will focus on...

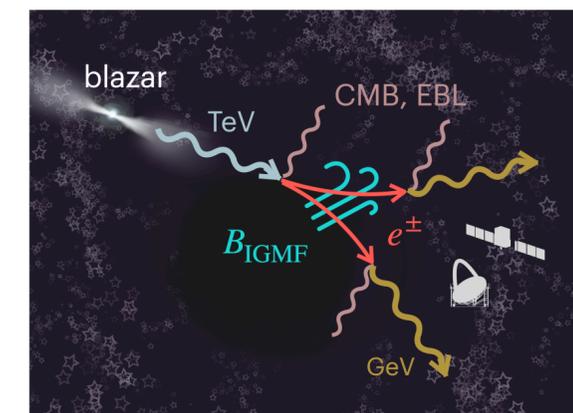
primordial physics



observations



[Chyży, Buta]



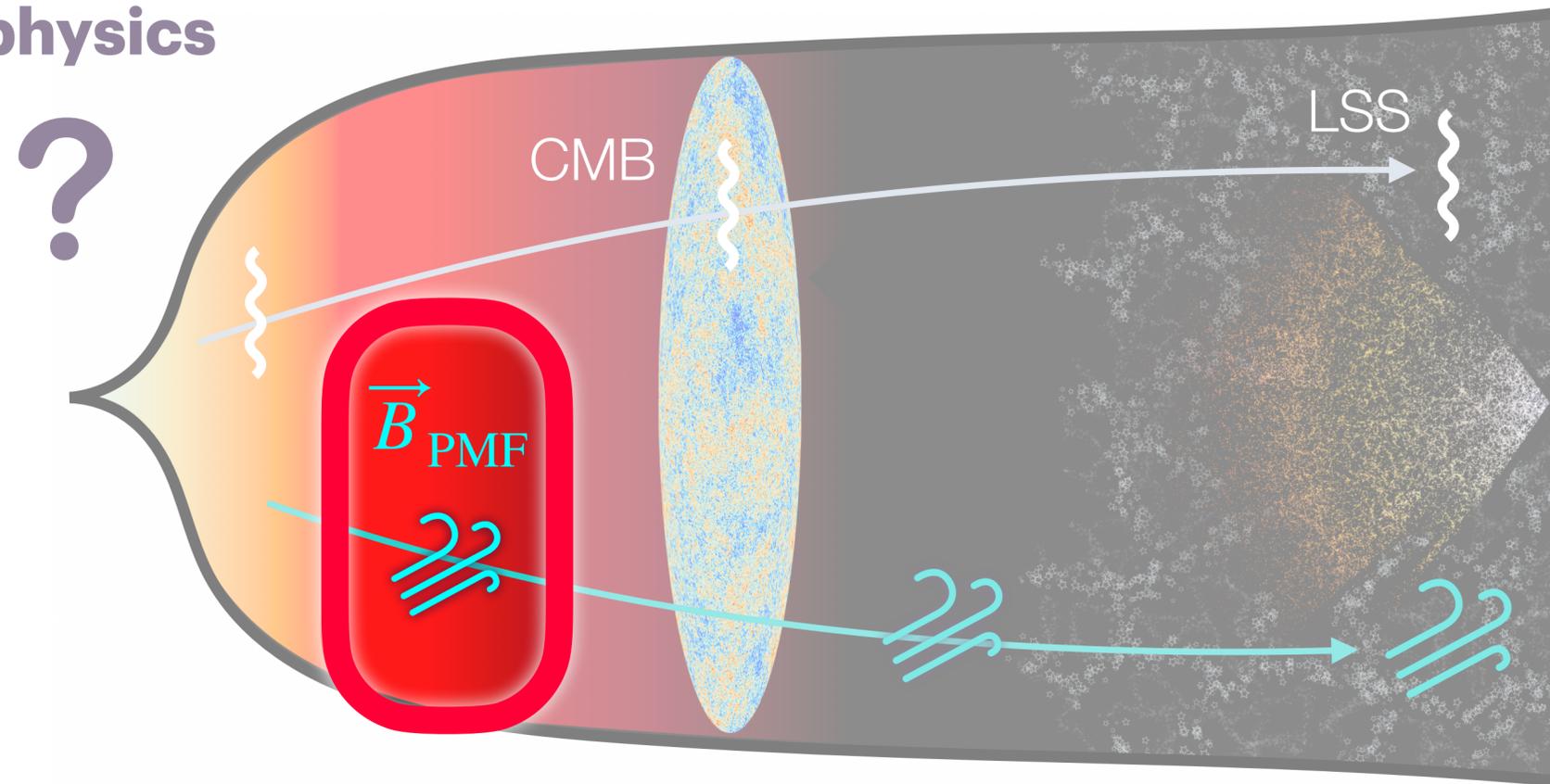
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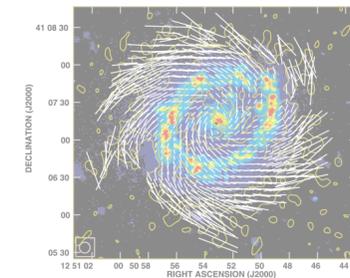
Figs (modified) from [Planck, ESA] and [D. Schlegel/Berkeley Lab using data from DESI, M. Zamani (NSF's NOIRLab)]

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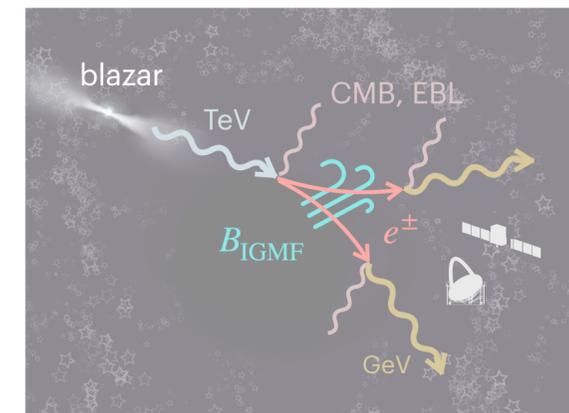
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[Chyży, Buta]



$B_{\text{void}} \gtrsim 10^{-17} \text{ G}$
[Neronov, Vovk 2010], ...

thermal plasma (SM) + magnetic field at $T \sim 100 \text{ GeV}$

- What “magnetic field” means is nontrivial
- Possibly relevant to baryogenesis

Figs (modified) from [Planck, ESA] and [D. Schlegel/Berkeley Lab using data from DESI, M. Zamani (NSF's NOIRLab)]

Outline

- How to describe “magnetic field” at ~ 100 GeV?
- Baryogenesis by primordial magnetic field?
- Can primordial magnetic field explain the void magnetic field?

Magnetic fields at high/low temperatures

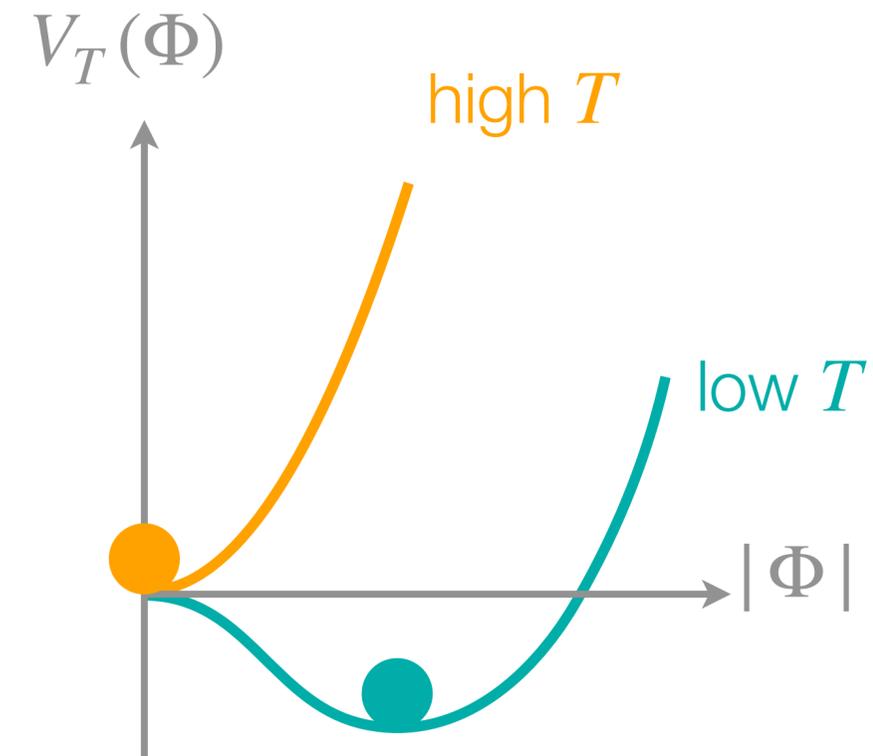
$$\mathcal{L}_{\text{SM}} \supset -\frac{1}{4}W_{\mu\nu}^a W^{a\mu\nu} -\frac{1}{4}Y_{\mu\nu} Y^{\mu\nu} + (D_\mu\Phi)^\dagger D^\mu\Phi + V(\Phi)$$

$$T \gg 100 \text{ GeV} \quad \langle |\Phi|^2 \rangle = 0$$

$U(1)_Y$ **massless field**
and massive W^1, W^2, W^3

$$T \ll 100 \text{ GeV} \quad \langle |\Phi|^2 \rangle \neq 0$$

$U(1)_{\text{em}}$ **massless field**
and massive W^\pm, Z



Transition during the electroweak crossover

$T \gg 100 \text{ GeV}$



$$\vec{B}_Y = \vec{\nabla} \times \vec{Y}$$

hyper-magnetic field



$T \ll 100 \text{ GeV}$



$$\vec{B}_{\text{em}} = \vec{\nabla} \times \vec{A}, \quad A_i = \cos \theta_w Y_i - \sin \theta_w W_i^3$$

the ordinary magnetic field

Naive interpolation

[D'Onofrio, Rummukainen 2015] [Kamada, Long 2016]

$T \gg 100 \text{ GeV}$



$$\theta_{\text{eff}}(T \gg 100 \text{ GeV}) = 0$$

$T \sim 100 \text{ GeV}$



$$\theta_{\text{eff}}(T)$$

effective mixing angle

$T \ll 100 \text{ GeV}$



$$\theta_{\text{eff}}(T \ll 100 \text{ GeV}) = \theta_w$$

Determining the effective mixing angle

$$T \ll 100 \text{ GeV} \quad \vec{B}_{\text{em}} = \cos \theta_w \vec{B}_Y - \sin \theta_w \vec{B}_{W^3}$$

$T \sim 100 \text{ GeV}$

effective mixing angle
 [D'Onofrio, Rummukainen 2015]
 [Kamada, Long 2016]

$$\cos \theta_{\text{eff}}(T) = \left(1 + \frac{g_3(T)^2 \sin^2 \theta_w}{4\pi m_W(T)} \right) \cos \theta_w$$

$$\begin{pmatrix} \vec{B}_\ell \\ \vec{B}_c \end{pmatrix} = \begin{pmatrix} \cos \theta_{\text{eff}}(T) & -\sin \theta_{\text{eff}}(T) \\ \sin \theta_{\text{eff}}(T) & \cos \theta_{\text{eff}}(T) \end{pmatrix} \begin{pmatrix} \vec{B}_Y \\ \vec{B}_{\mathcal{W}} \end{pmatrix}$$

gauge invariance $\mathcal{W}_{ij} := -\frac{\Phi^\dagger \sigma^a \Phi}{\Phi^\dagger \Phi} W_{ij}^a$

$$\text{s.t. } \langle B_{\ell i}(\vec{p}) B_{\ell j}(\vec{q}) \rangle \sim Z^\ell(T) P_{ij}(\hat{p}) + \dots, \dots$$

massless pole
 [Kajantie+ 1997]

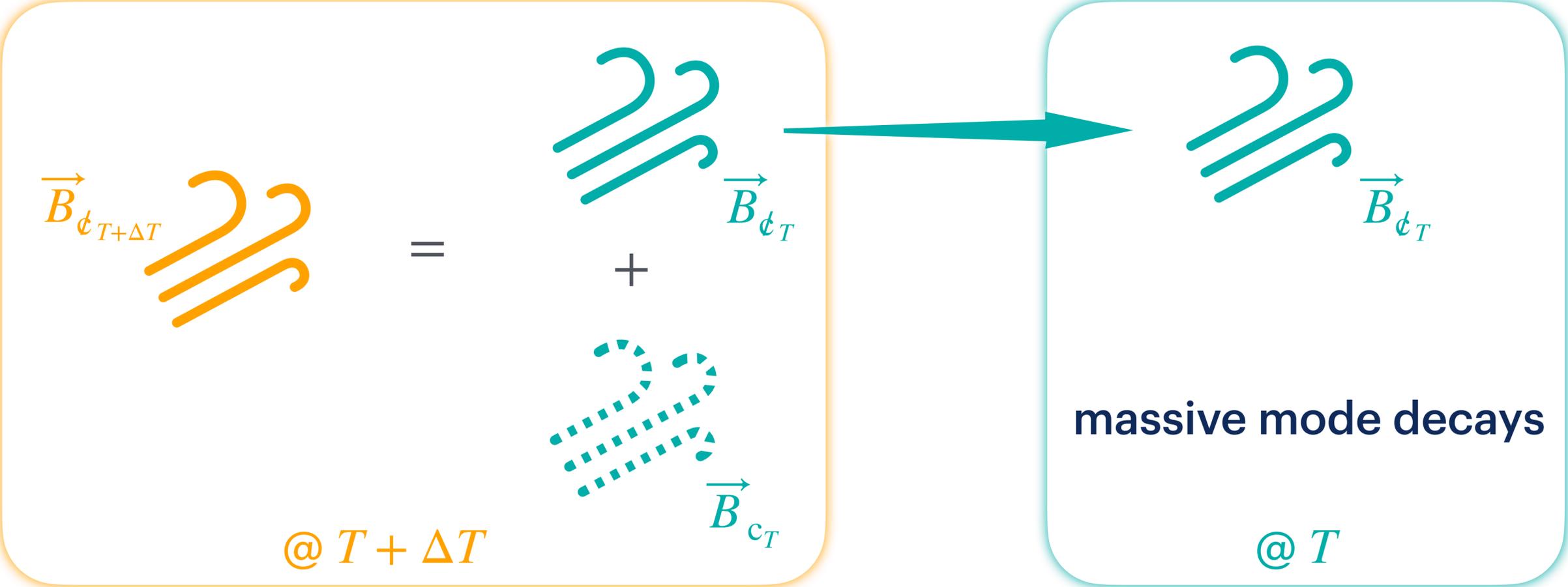
wavefunction renormalization
 [Hamada, Mukaida, FU 2025a]

$$Z^\ell(T) = 1 + \left(\frac{5}{12} - \frac{1}{24 \sqrt{1 + \frac{m_D(T)^2}{m_W(T)^2}}} \right) \frac{g_3(T)^2 \sin^2 \theta_w}{\pi m_W(T)}$$

Intuitive picture

In terms of massless $\not\epsilon_T$ and massive c_T modes at every moment T ,

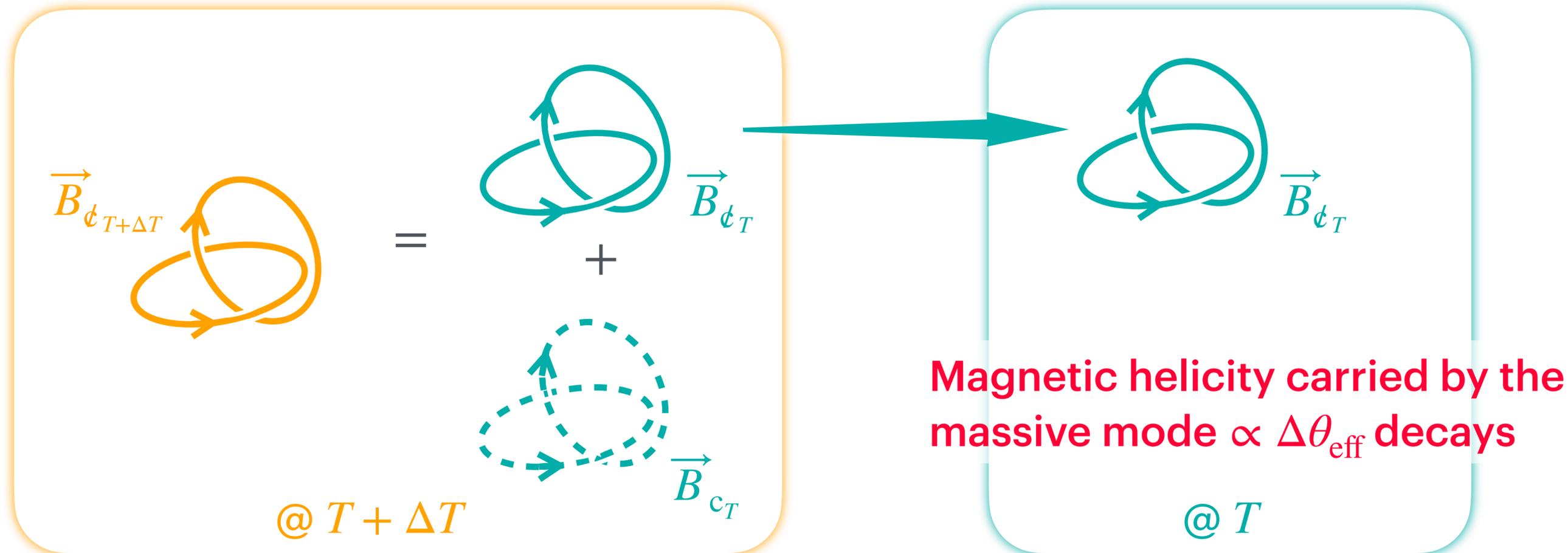
$$\vec{B}_{\not\epsilon_{T+\Delta T}} = \cos \Delta\theta_{\text{eff}} \vec{B}_{\not\epsilon_T} + \sin \Delta\theta_{\text{eff}} \vec{B}_{c_T}, \quad \Delta\theta_{\text{eff}} := \theta_{\text{eff}}(T + \Delta T) - \theta_{\text{eff}}(T)$$



Helicity decay $\propto \Delta\theta_{\text{eff}}$

magnetic helicity $\int d^3x \vec{A}_{\mathfrak{c}} \cdot \vec{B}_{\mathfrak{c}} =$ topological charge of magnetic field lines [Moffatt 1969]

$$\vec{B}_{\mathfrak{c}_{T+\Delta T}} = \cos \Delta\theta_{\text{eff}} \vec{B}_{\mathfrak{c}_T} + \sin \Delta\theta_{\text{eff}} \vec{B}_{\mathfrak{c}_T}, \quad \Delta\theta_{\text{eff}} := \theta_{\text{eff}}(T + \Delta T) - \theta_{\text{eff}}(T)$$



Anomaly equation roughly says...

[Joyce, Shaposhnikov 1997] [Giovannini, Shaposhnikov 1998] [Fujita, Kamada 2016]...

[Kamada, Long 2016]

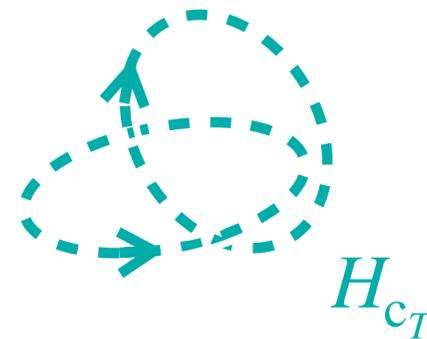
$$\Delta(\text{baryon asymmetry}) \sim - \Delta(\text{magnetic helicity})$$

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$$B > \bar{B}$$



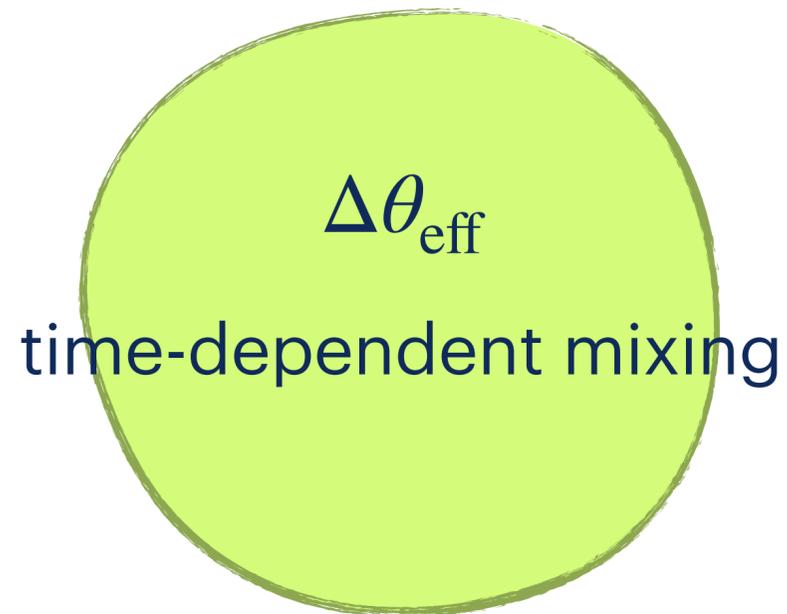
More precisely, ...

[Joyce, Shaposhnikov 1997] [Giovannini, Shaposhnikov 1998] [Fujita, Kamada 2016]...

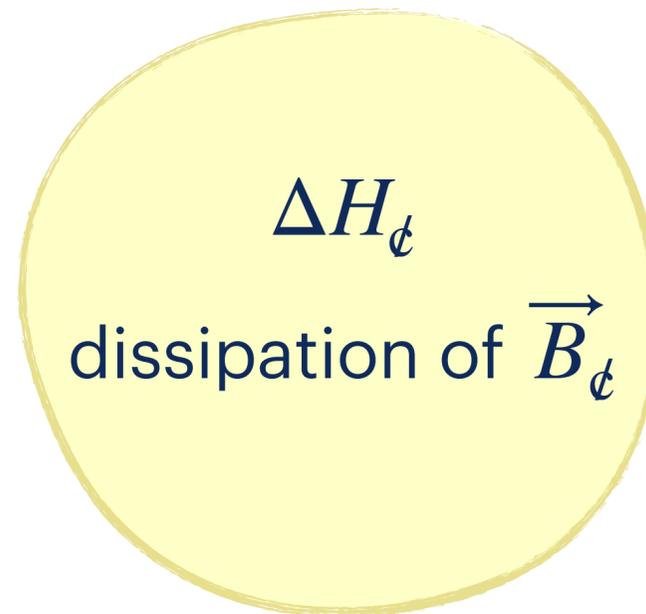
[Kamada, Long 2016] [Hamada, Mukaida, FU 2025b]

chiral anomaly [Adler 1969] [Bell, Jackiw 1969]

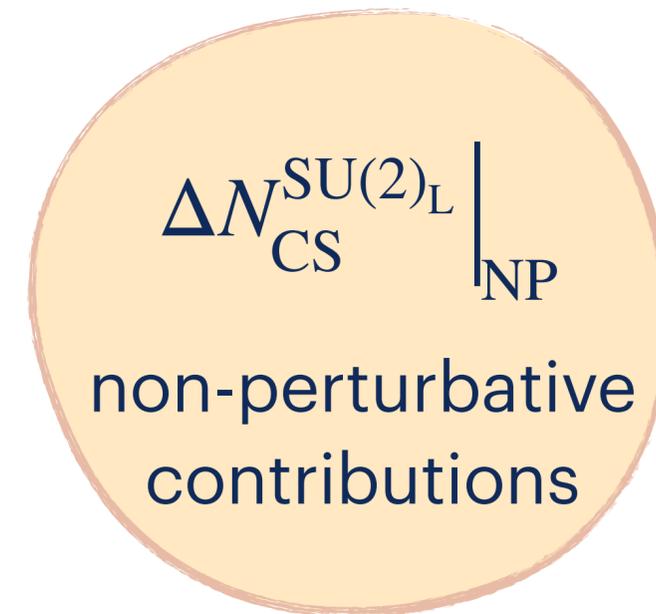
$$\begin{aligned}\Delta Q_{B+L} &= 2 \cdot 3 \left(\Delta N_{CS}^{\text{SU}(2)_L} - \Delta H_Y \right) \\ &\simeq 3 \cdot 2 \cot^2 \theta_w \Delta \left[\left(\tan^2 \theta_{\text{eff}} - \tan^2 \theta_w \right) H_\phi \right] + 3 \cdot 2 \Delta N_{CS}^{\text{SU}(2)_L} \Big|_{\text{NP}}\end{aligned}$$



helicity decay



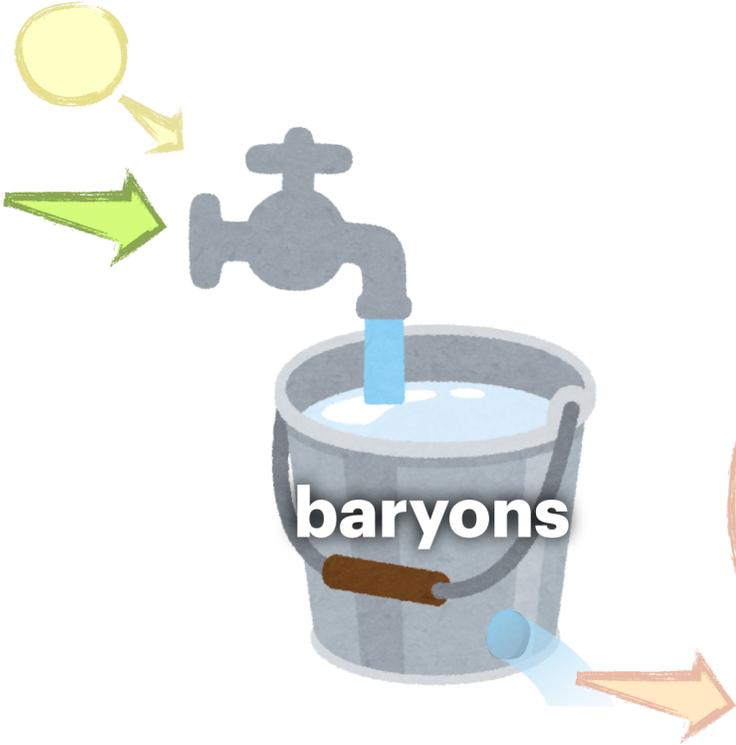
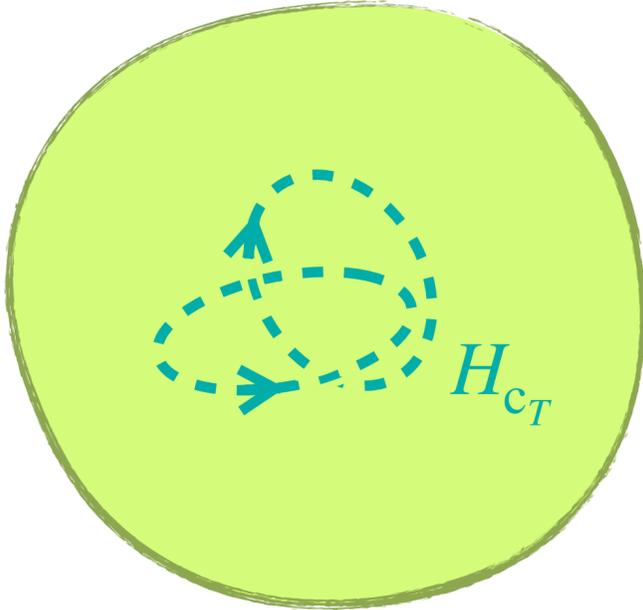
MHD



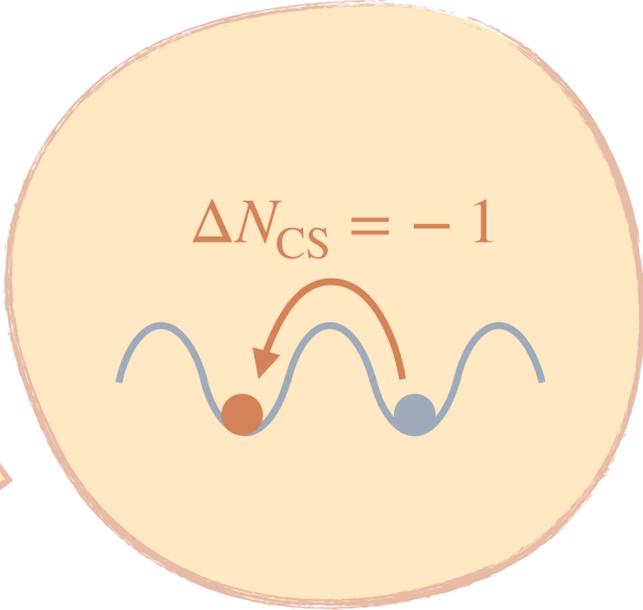
sphaleron

Baryogenesis from magnetic helicity decay

magnetic helicity decay



sphaleron washout



equilibrium until ~ 130 GeV
[Kamada, Long 2016]

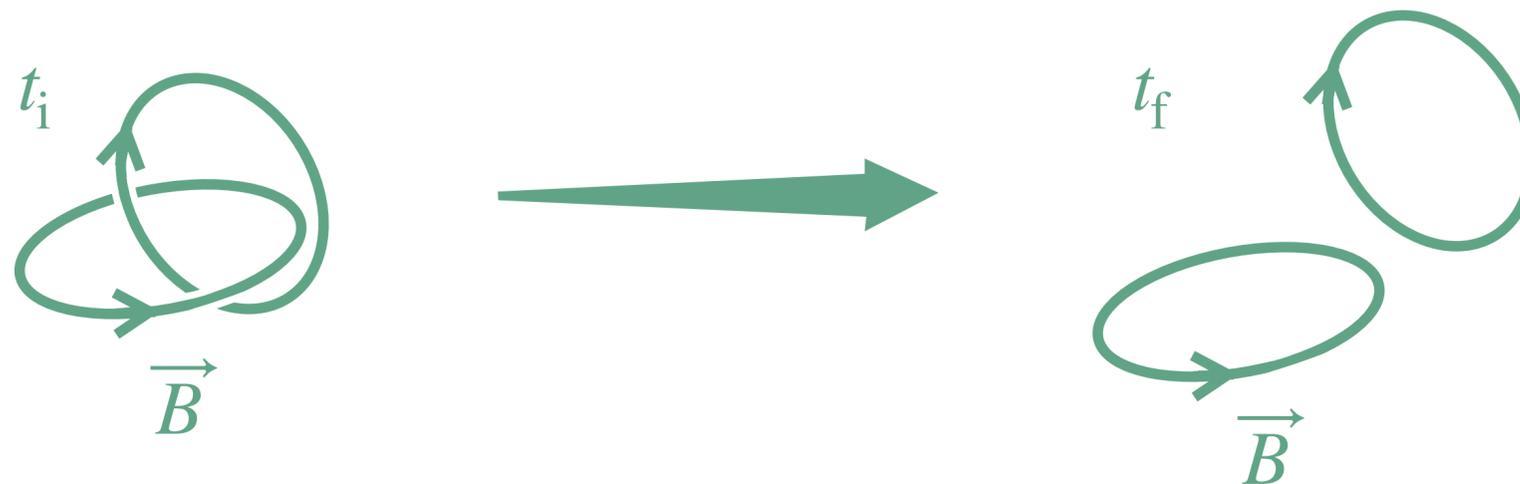
the baryon yield can be examined under BBN

..., [Fujita, Kamada 2016] [Kamada, Long 2016]
[Giovannini, Shaposhnikov 1997] [Kamada, FU, Yokoyama, 2021]

However, we may have missed another washout effect.

Magnetic helicity can be ill-defined

Assume $T = 0$ and $\theta_{\text{eff}} = \theta_w$ for simplicity.

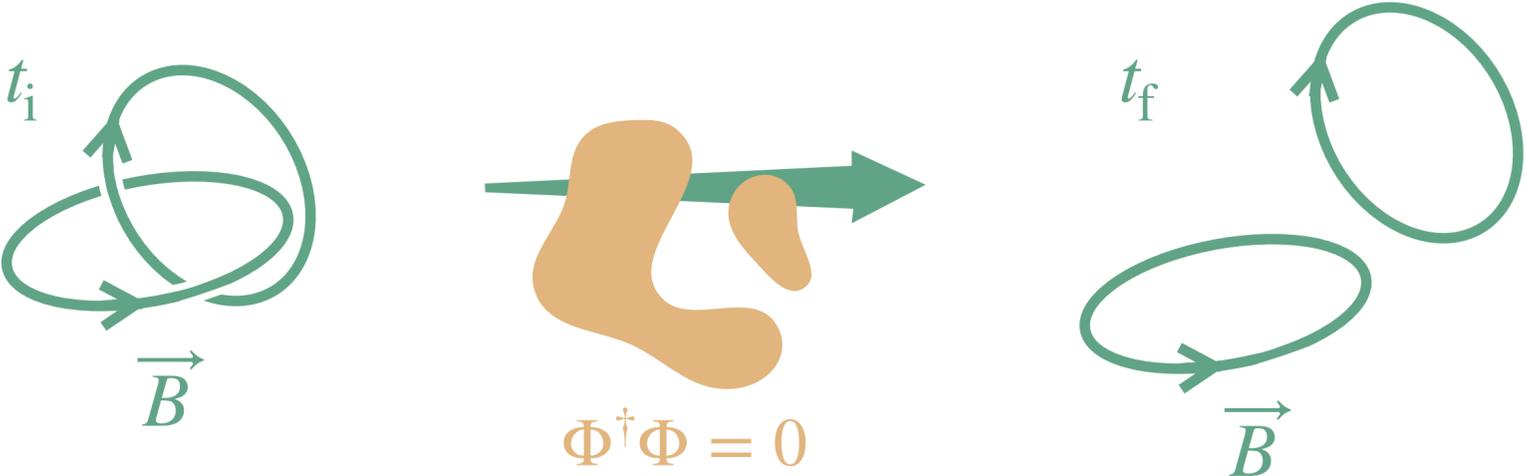


Naively,
$$\Delta \int d^3x \vec{A} \cdot \vec{B} \stackrel{?}{=} -2 \int_{t_i}^{t_f} d^4x \vec{E} \cdot \vec{B}$$

Magnetic helicity can be ill-defined

Assume $T = 0$ and $\theta_{\text{eff}} = \theta_w$ for simplicity.

During $t_i < t < t_f$, the “ $U(1)_{\text{em}}$ ” can be ill-defined in the region where $\Phi^\dagger \Phi = 0$



$$\Delta \int d^3x \vec{A} \cdot \vec{B} = -2 \int_{t_i}^{t_f} d^4x \vec{E} \cdot \vec{B} + \int_{\partial \mathcal{L}} A \wedge F \quad \text{nontrivial helicity change where higgs vanishes!}$$

[Fukuda, Hamada, Kamada, Mukaida, FU 2025]

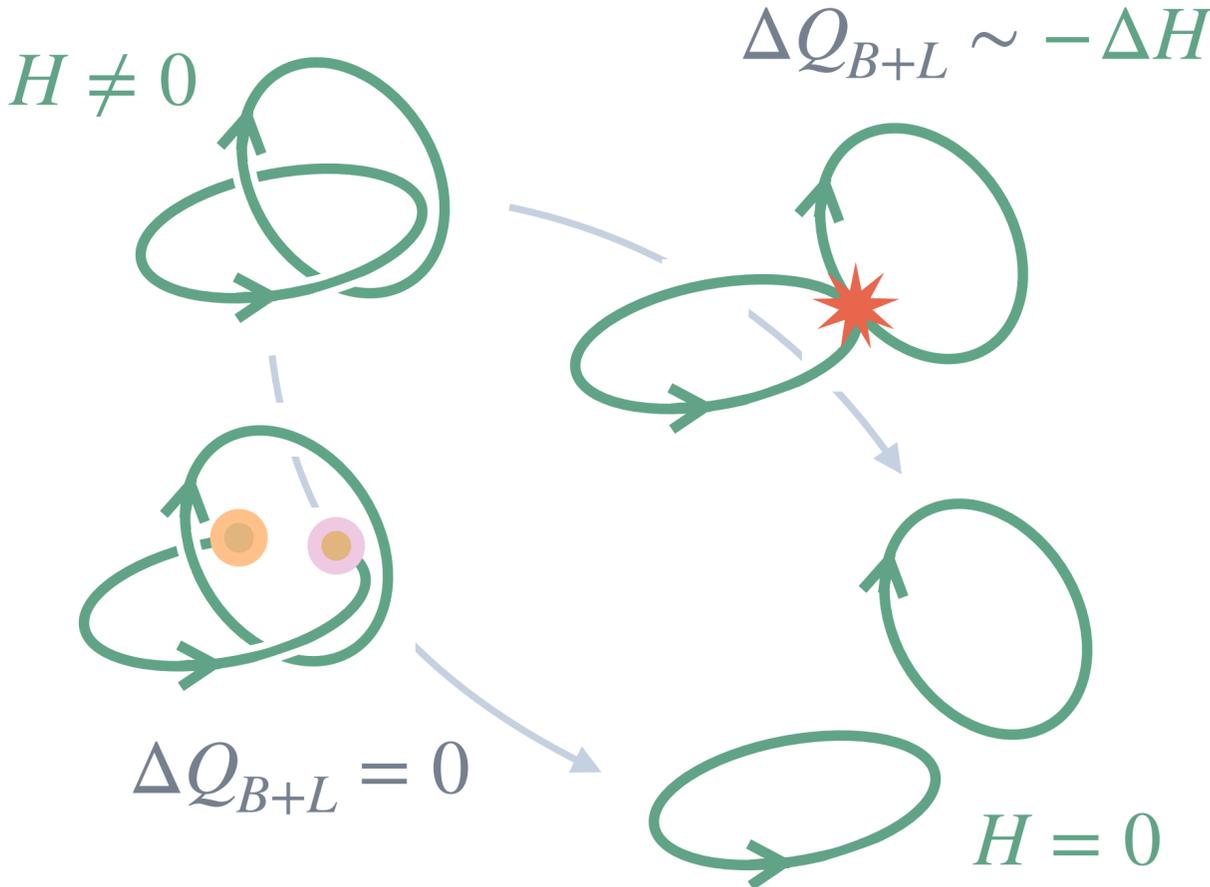
Indeed, Higgs dynamics can matter

Toy model (non-SM)

$$SU(2) \rightarrow U(1) \rightarrow 1, \quad \partial_\mu j_{B+L}^\mu \sim \vec{E} \cdot \vec{B}$$

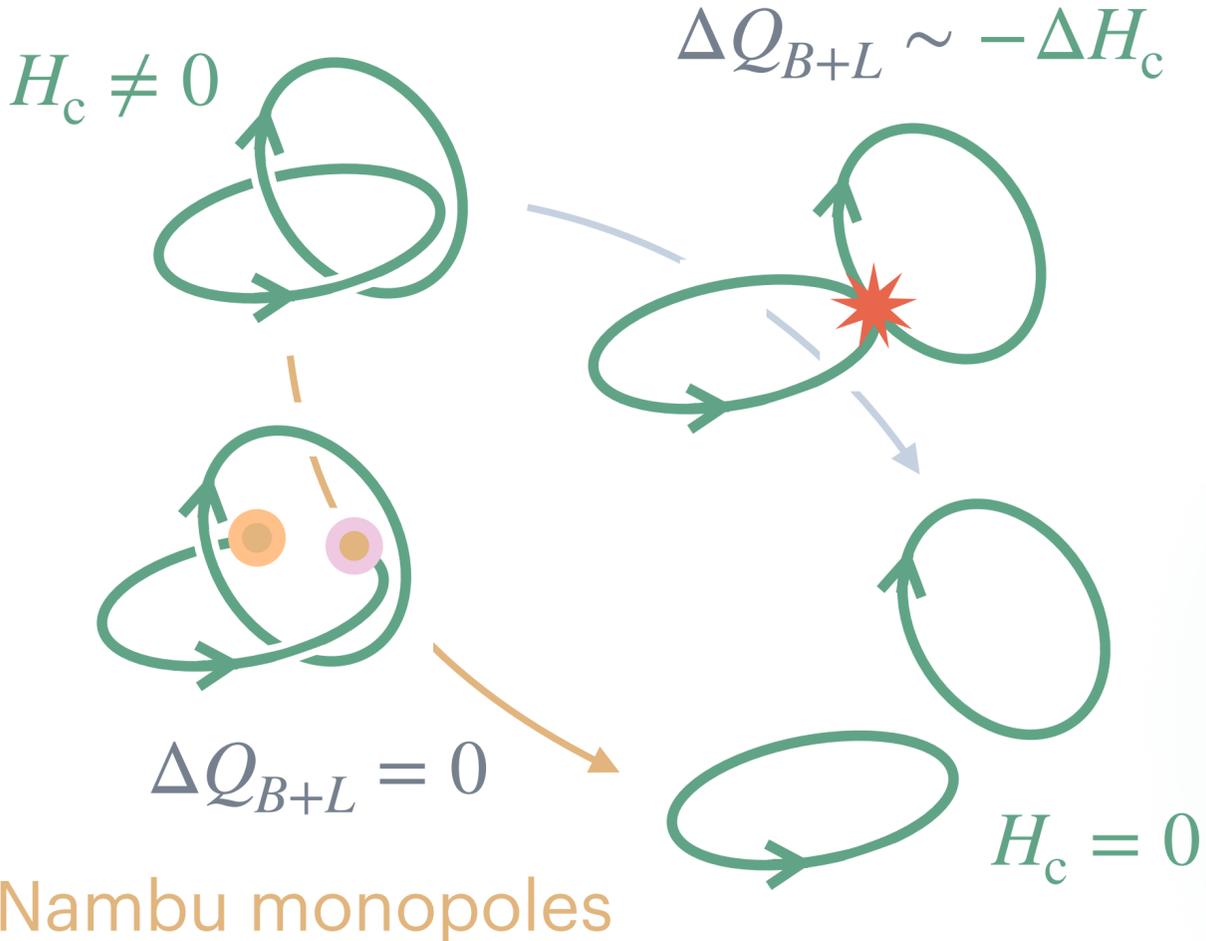
$\Delta Q_{B+L} = -\Delta H$ does not hold always.

untying link by generating
't Hooft—Polyakov monopole pairs

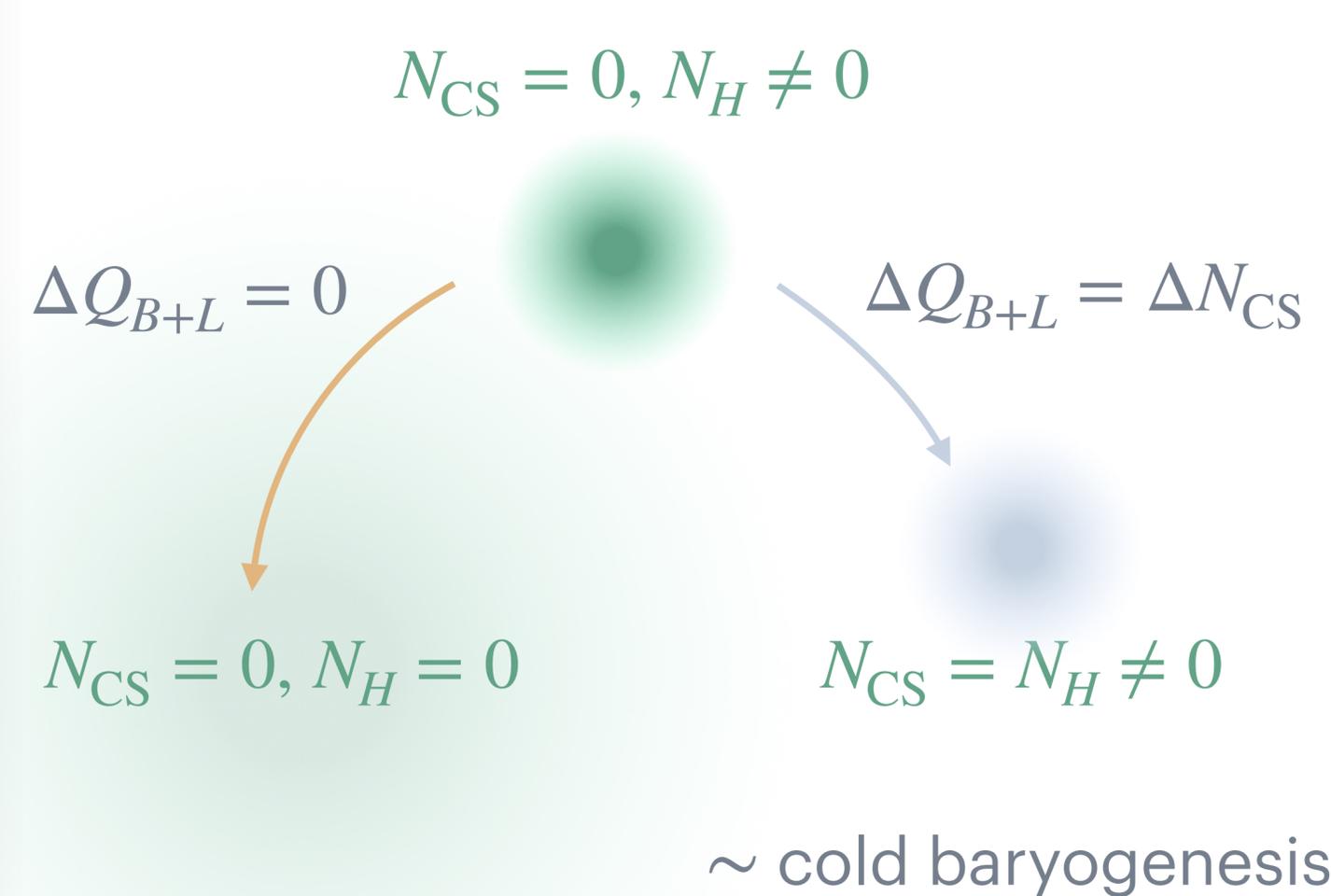


In the SM? : future work

Untying links?



Texture decay?



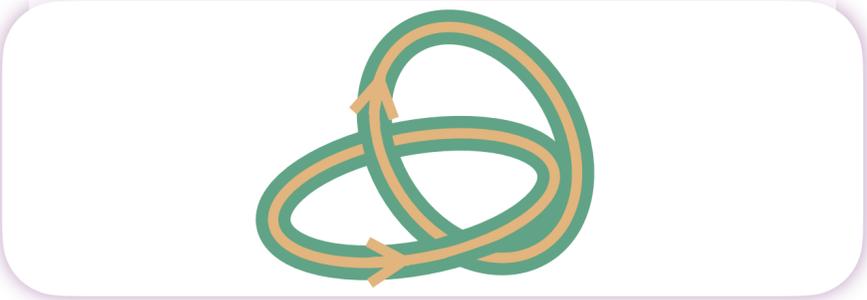
Putting it in another way, ...



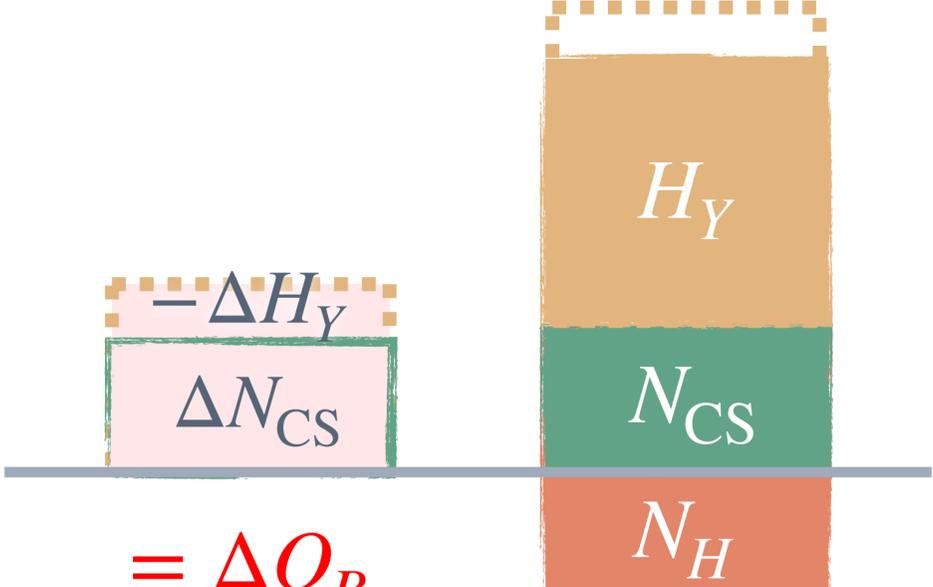
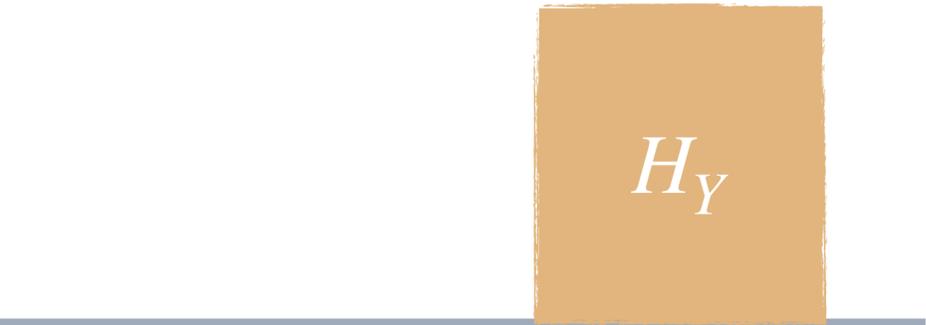
@ $T \gg 100 \text{ GeV}$



$\vec{B}_{em} = \vec{B}_Y$ dressed by \vec{B}_W



@ $T \ll 100 \text{ GeV}$



$= \Delta Q_B$

$$N_{CS} - H_Y - N_H = 0$$

Baryon asymmetry constraints [Hamada, Mukaida, FU 2025b]

helical $U(1)_Y$ magnetic field

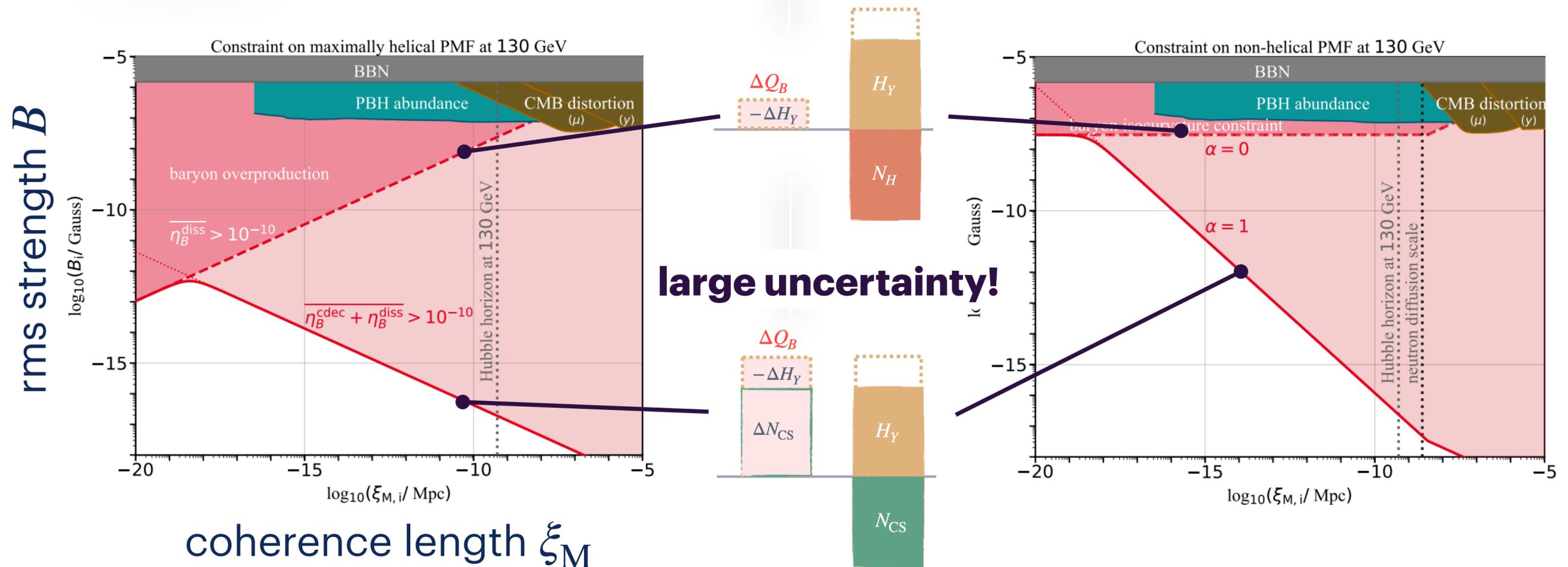
$$\langle \delta\eta_B \rangle \sim \langle \vec{A} \cdot \vec{B} \rangle$$

$$\leq 10^{-10} \text{ (CMB, BBN)}$$

non-helical $U(1)_Y$ magnetic field

$$\langle \delta\eta_B^2 \rangle \sim \langle (\vec{A} \cdot \vec{B})^2 \rangle, \quad \langle \vec{A} \cdot \vec{B} \rangle = 0$$

$$\leq \mathcal{O}(0.01) \eta_{B,\text{obs}}^2 \text{ (BBN) [Inomata+ 2018]}$$



Baryon asymmetry constraints [Hamada, Mukaida, FU 2025b]

helical $U(1)_Y$ magnetic field

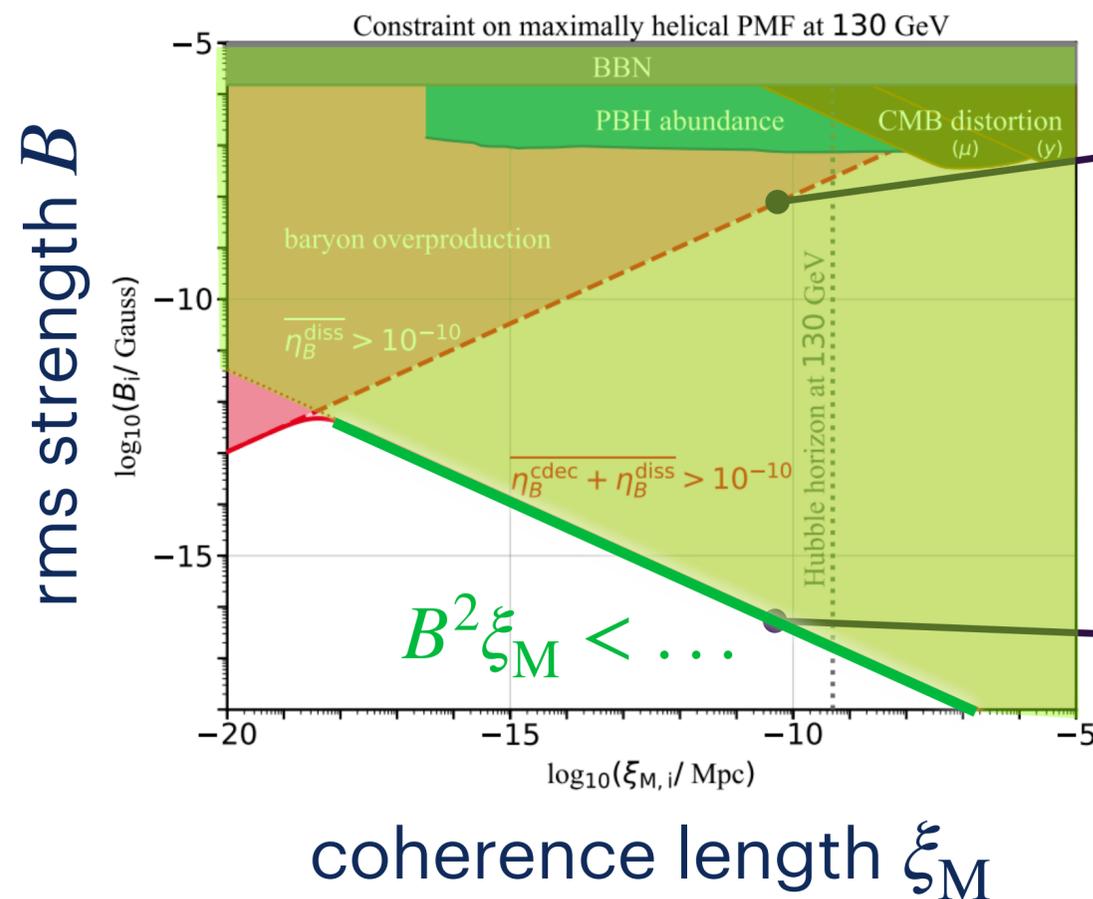
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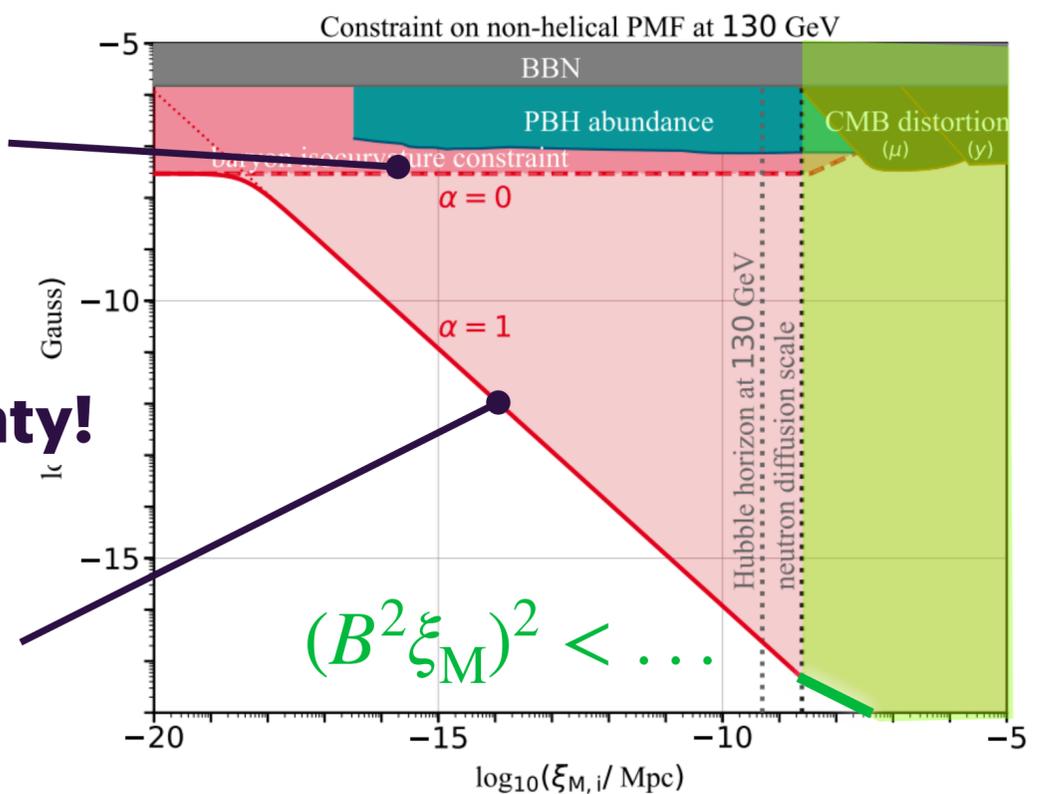
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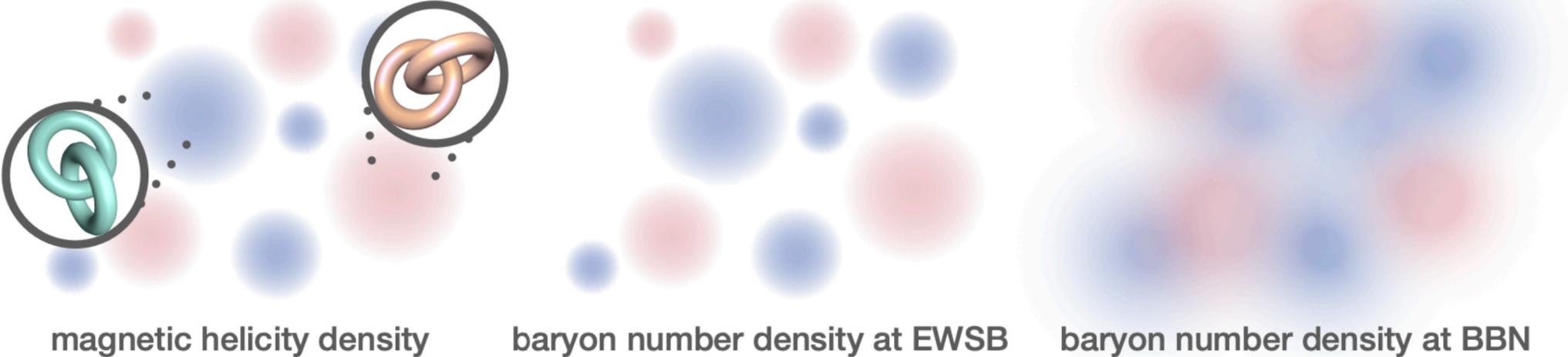
large uncertainty!



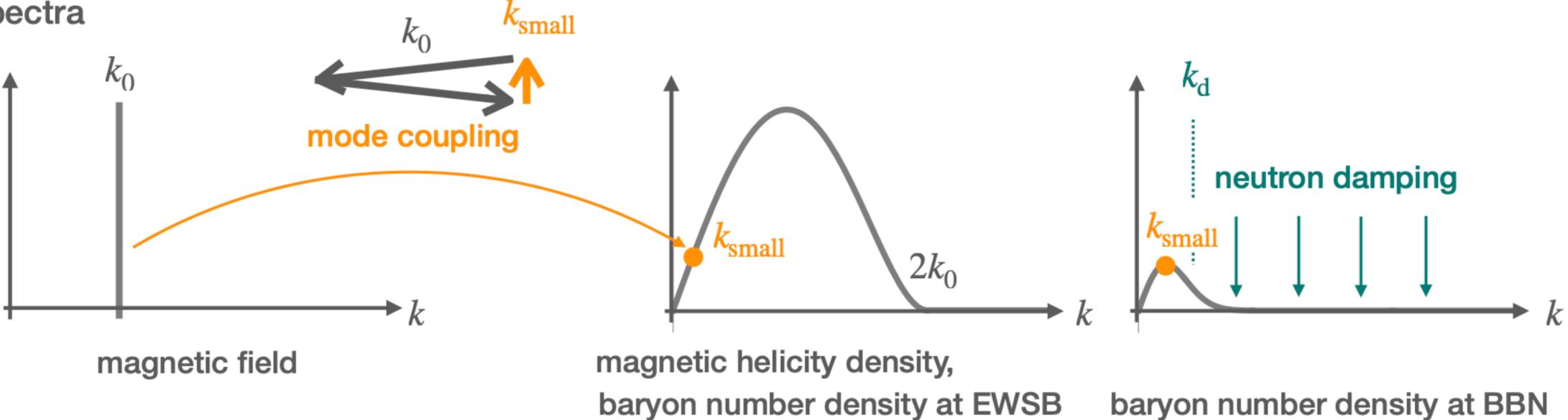
Baryon perturbation at small scales

[Kamada, FU, Yokoyama 2021] [Akiyama, FU ongoing]

real space configurations

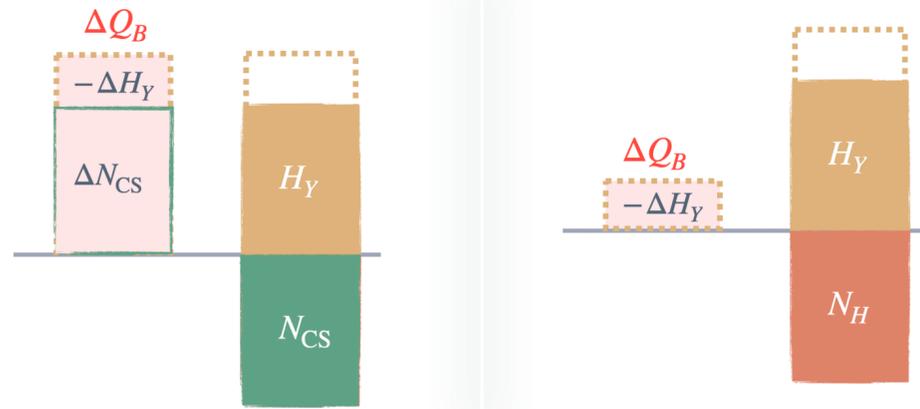


power spectra



Originally- $U(1)_Y$ void magnetic field?

the constraints + **MHD evolution**



If $\Delta N_H = 0$ / $\Delta N_{CS} \simeq 0$

○ at best



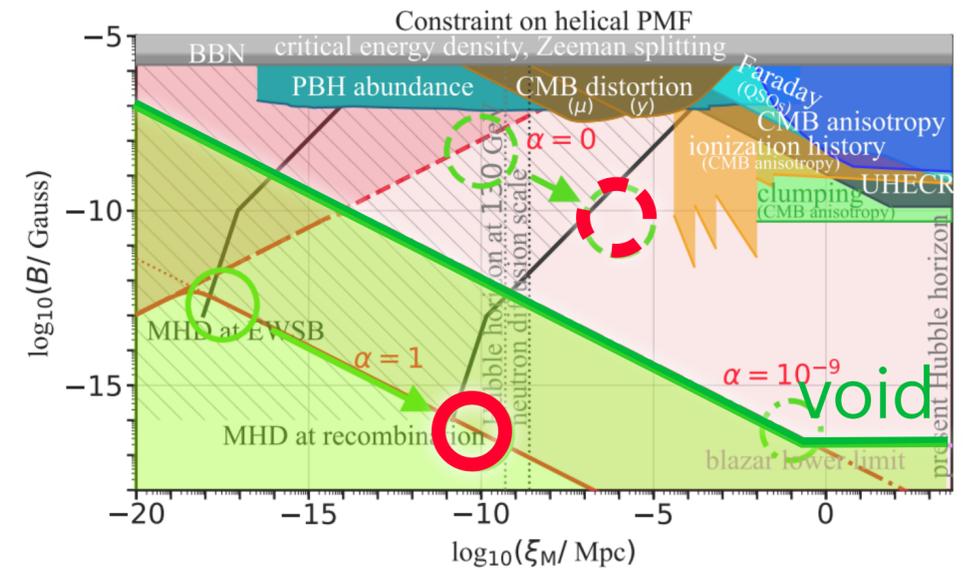
PMF $\rightarrow \vec{B}_{\text{void}}$



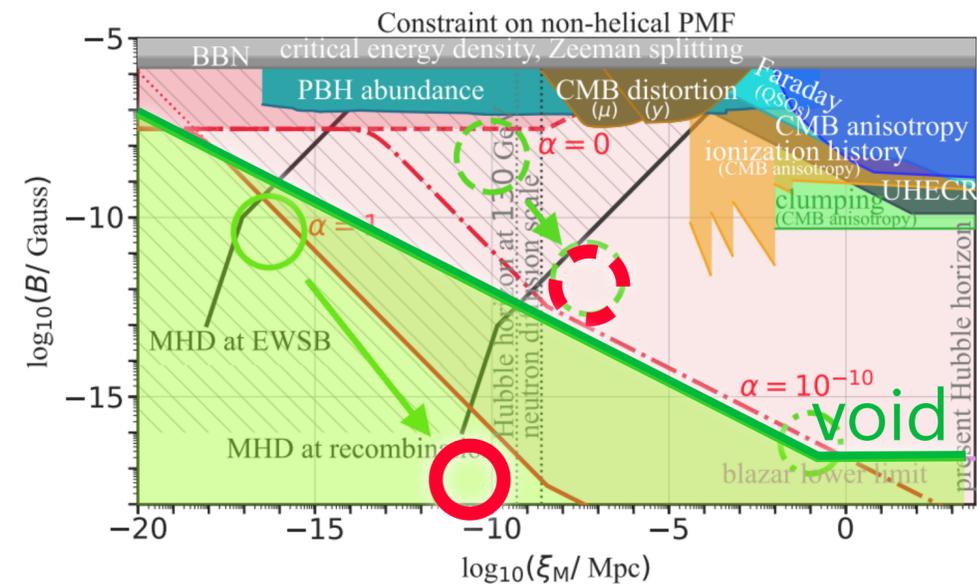
PMF $\rightarrow \eta_{B, \text{obs}}$



baryon overproduction (helical)



baryon fluctuation (non-helical)



[Hamada, Mukaida, FU 2025b]

Summary

- **continuous transition from \vec{B}_Y to \vec{B}_{em}** during the electroweak crossover
 - ✓ one-loop formula for effective mixing angle $\theta_{\text{eff}}(T)$ and normalization $Z^{\mathcal{C}}(T)$
- **Helicity decay** $\propto \Delta\theta_{\text{eff}}$ at every instant
 - ✓ may **generate baryons** through the chiral anomaly, (known in the literature,)
 - ✓ but **we have missed possible contribution of the Higgs dynamics.**
- Once **ΔN_H during the crossover** is known (future work), one can estimate ΔQ_B .
 - ✓ If $\Delta N_H = 0$, generated baryon asymmetries suffer from BBN constraints.
 - ✓ Only if $\Delta N_H \neq 0$ s.t. $\Delta N_{CS} \simeq 0$, **\vec{B}_Y may be origins of both \vec{B}_{void} and $\eta_{B,\text{obs}}$.**